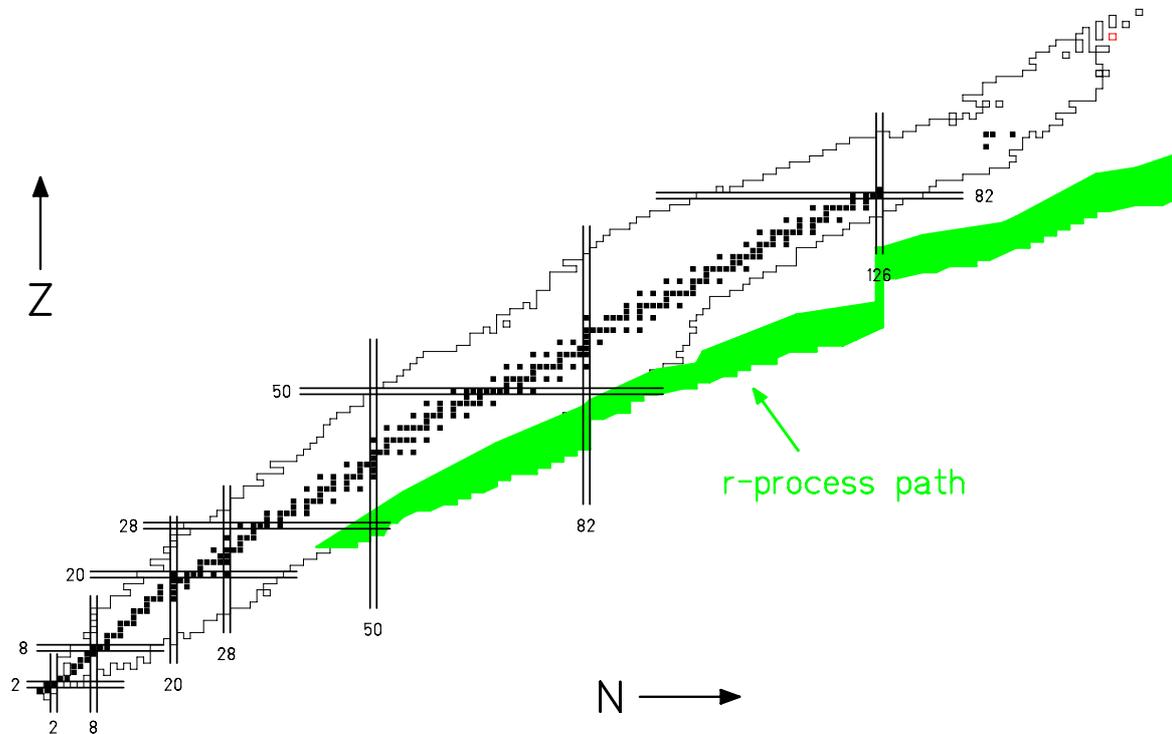


Prospects for the production of neutron-rich nuclei

(Karl-Heinz Schmidt, GSI Darmstadt)

- 1. Arguments for an optimum beam energy**
- 2. Nuclear reactions below the Fermi energy**
- 3. Nuclear reactions above the Fermi energy**
- 4. Fusion**
- 5. Evaporation**
- 6. Fission**
- 7. Fragmentation**
- 8. Two-step reactions?**

Finding best conditions to go beyond the present limits

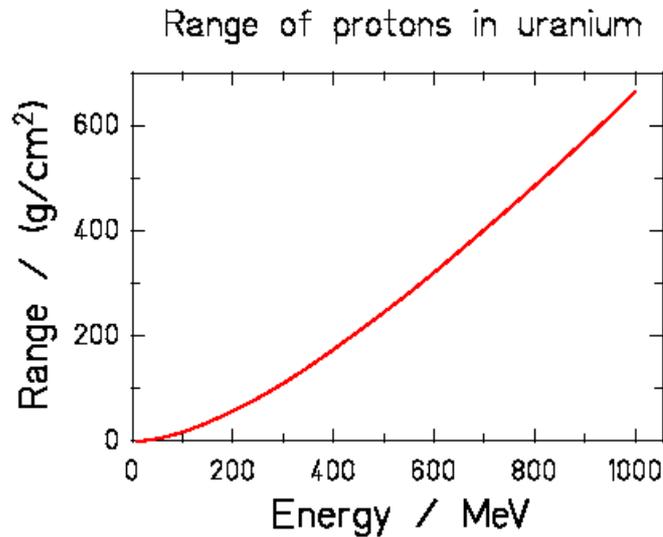


Even nature did not reach the neutron drip line (r-process)

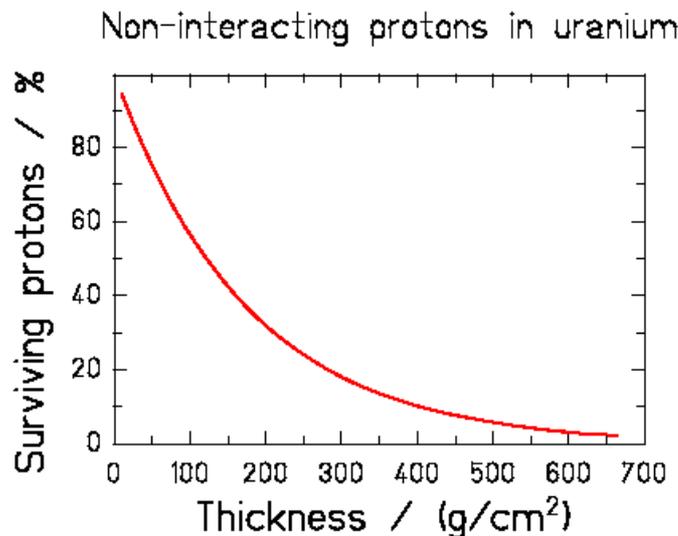
How to produce these neutron-rich nuclei in laboratory?

Range and usable target thickness

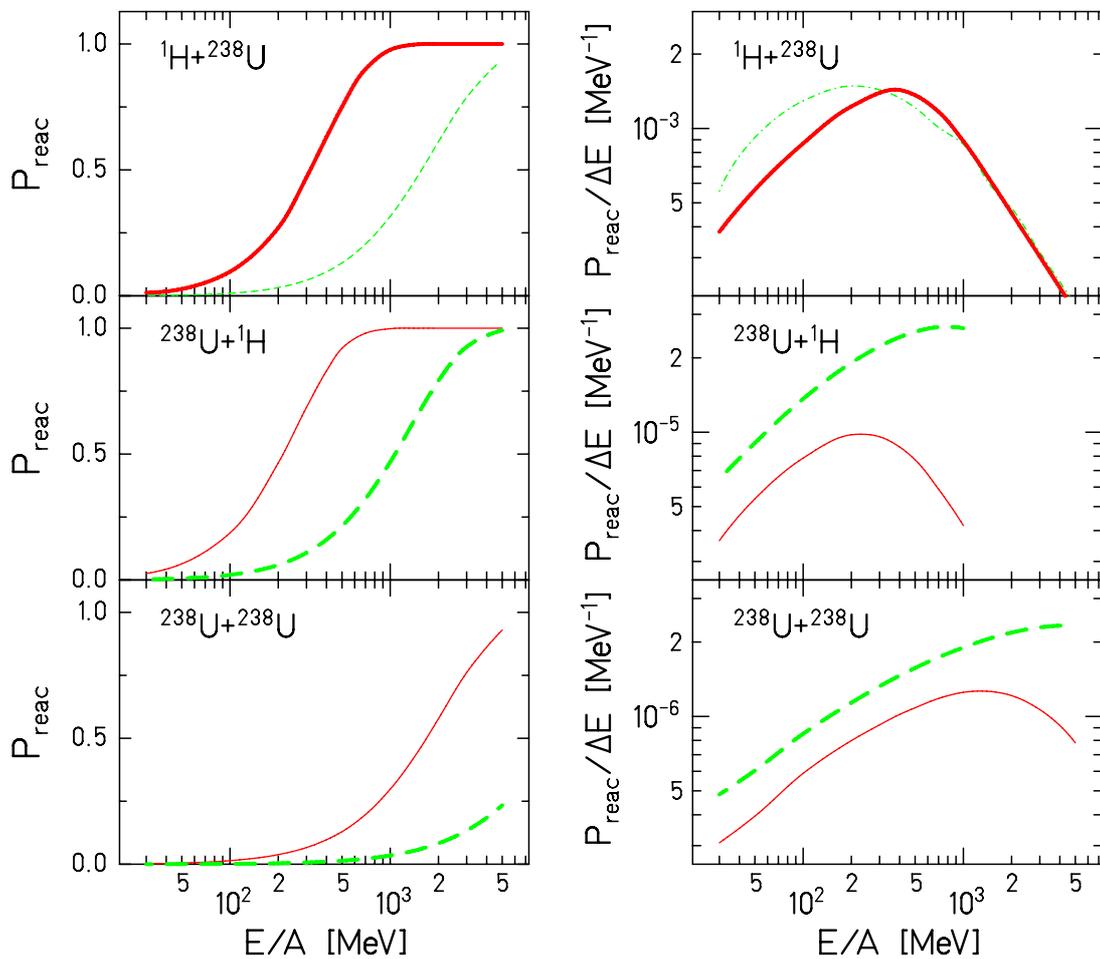
- *Range increases strongly with energy*



- *Usable target thickness limited by reaction probability*



General arguments for an "optimum" beam energy



P_{reac} = nuclear-reaction probability of projectile

ΔE = energy deposit in target per projectile

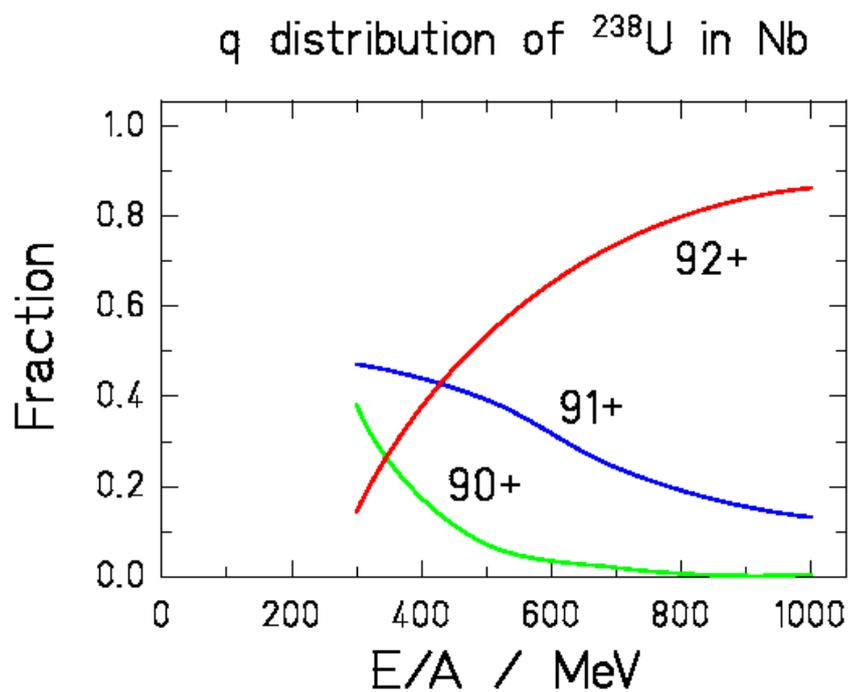
**Red lines: target thickness = range of projectiles
typical for ISOL-type facilities.**

**Green lines: target thickness = $0.1 \times$ range of projectiles
typical for in-flight facilities.**

**High reaction probabilities and low heat load
at $E \approx 1$ A GeV.**

Ionic charge states

- *Heavy reaction products need to be relativistic for in-flight separation*



Reaction mechanisms

$$E/A < E_{fermi}$$

Low reaction probabilities

Reactions controlled by nuclear potential and binding energies

Fusion, deep inelastic, transfer, fission

$$E/A > E_{fermi}$$

High reaction probabilities

Two stages:

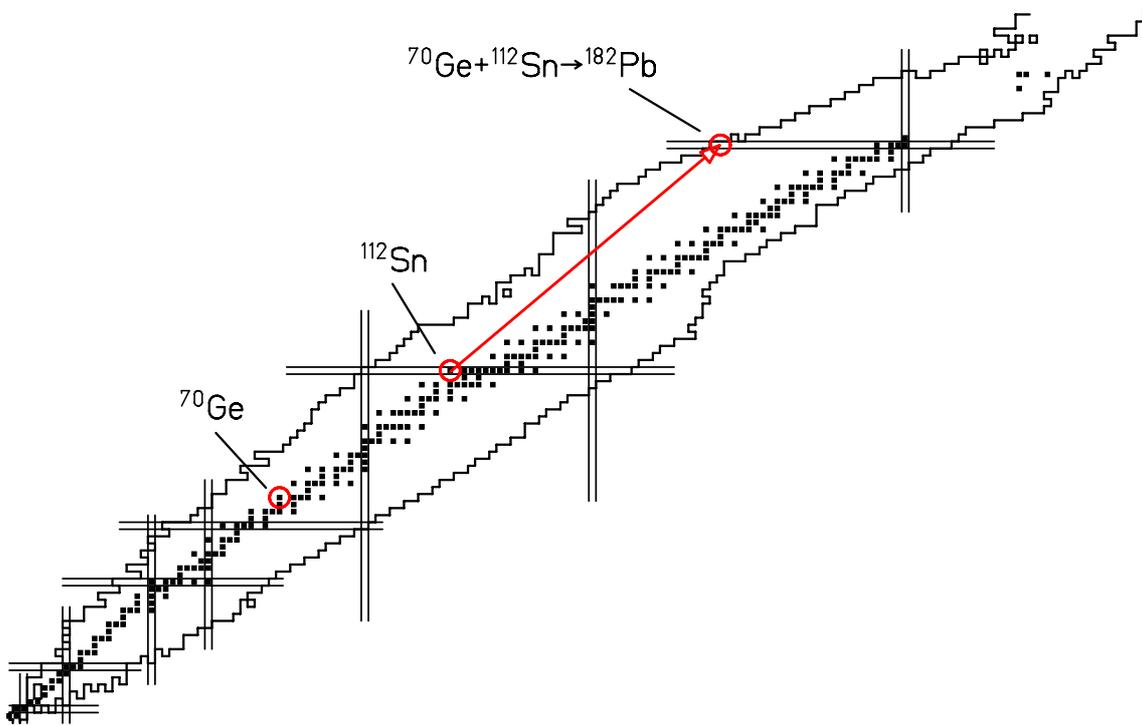
1. Collisions of individual nucleons

Target (resp. projectile) fragmentation

2. Deexcitation: nuclear potential and binding energies are again important

Deexcitation of prefragments by evaporation-fission competition

Fusion



- **Basic characteristics:**

Curvature of stability valley due to Coulomb repulsion.

→ **Fusion products are neutron-deficient.**

Nucleons of projectile and target add up.

→ **Suited to reach super-heavy nuclei.**

Partial widths Γ_n and Γ_p for emission of neutrons and protons.

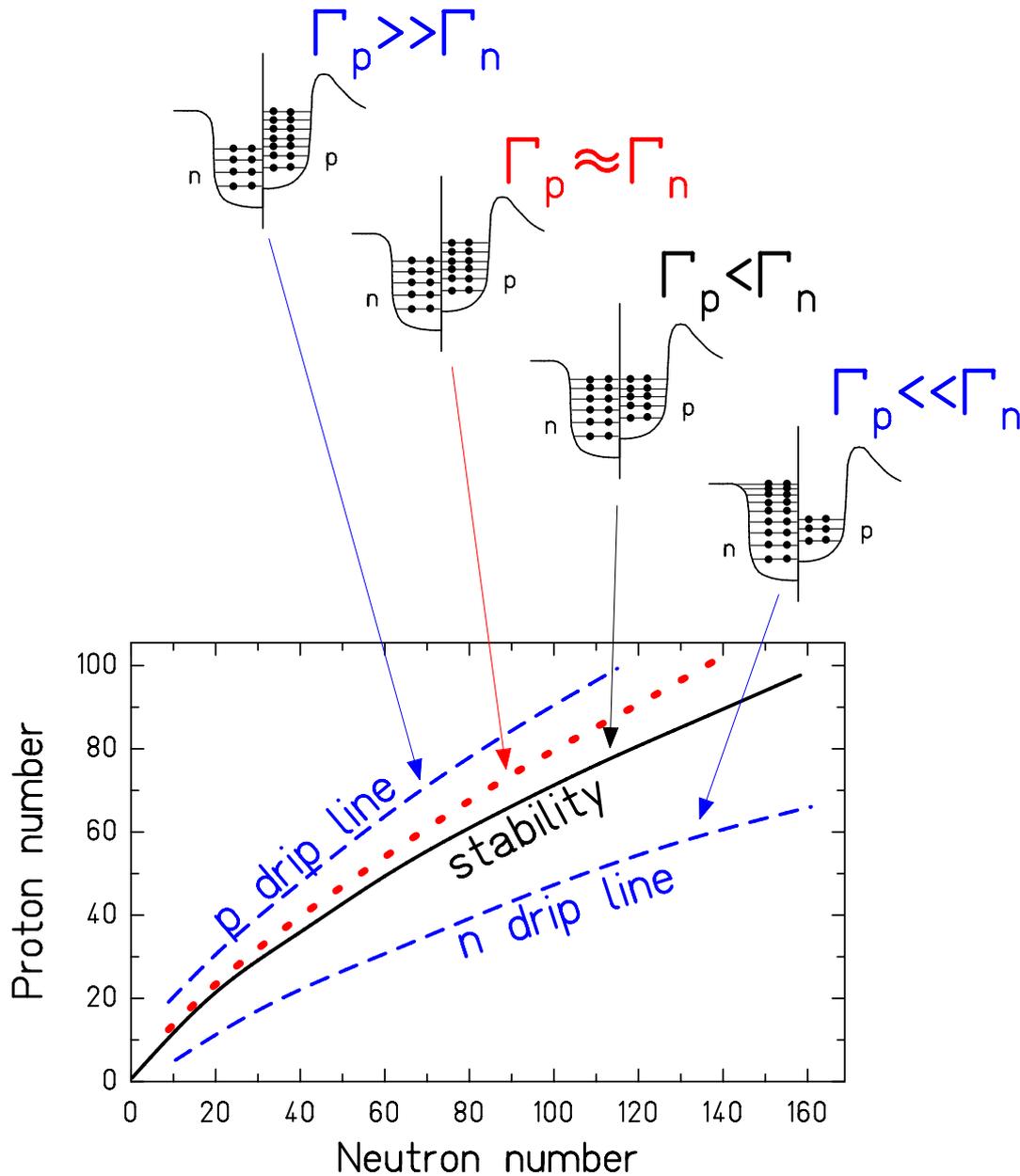
$$\Gamma_n = \frac{2mR^2 g}{2\pi\rho(E - E_r^{gs})} \int_0^{E-B_n} \varepsilon \rho(E - B_n - \varepsilon) d\varepsilon$$

$$\Gamma_p = \frac{2mR^2 g}{2\pi\rho(E - E_r^{gs})} \int_{\varepsilon_c}^{E-B_p} \varepsilon \left(1 - \frac{\varepsilon_c}{\varepsilon}\right) \rho(E - B_p - \varepsilon) d\varepsilon$$

(Approximation without considering tunneling.)

Γ_p is reduced by the Coulomb barrier ε_c .

Evaporation Process



Evaporation is like a diffusion process.

→ **Residues concentrate on an evaporation corridor.**

Proton evaporation is hindered by the Coulomb barrier.

→ **The evaporation corridor is neutron-deficient.**

Modelling the Width in A and N/Z of Fission-Product Isotopic Distributions

Approximated parabolic potential

$$U(\eta) = C_{\eta} \cdot (\eta - \eta_0)^2$$

Statistical population:

$$Y(\eta) \propto \exp\left\{2\sqrt{a(U_0 - U(\eta))}\right\} \quad \rightarrow$$

$$Y(\eta) \propto \exp\left\{-\frac{(\eta - \eta_0)^2}{2\sigma_{\eta}^2}\right\} \quad \text{with}$$

$$2\sigma_{\eta}^2 = \frac{T}{C_{\eta}}$$

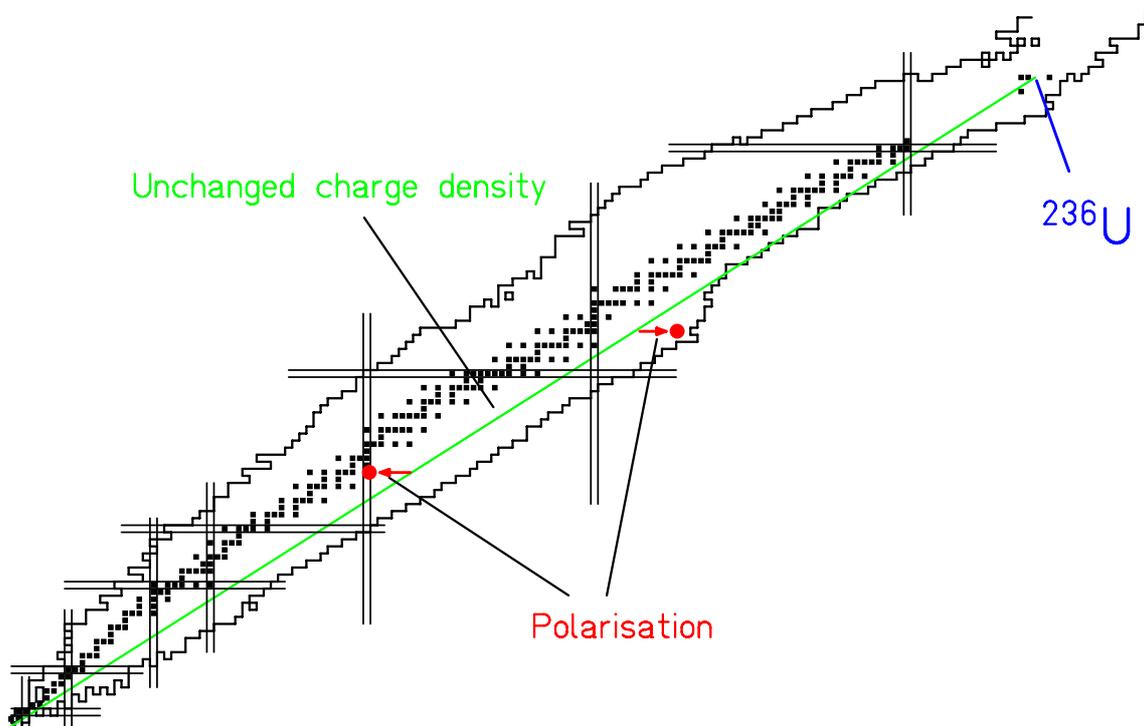
U = potential energy,

η = either A (mass split) or N/Z (polarisation),

C_{η} = stiffness of the potential,

T = nuclear temperature.

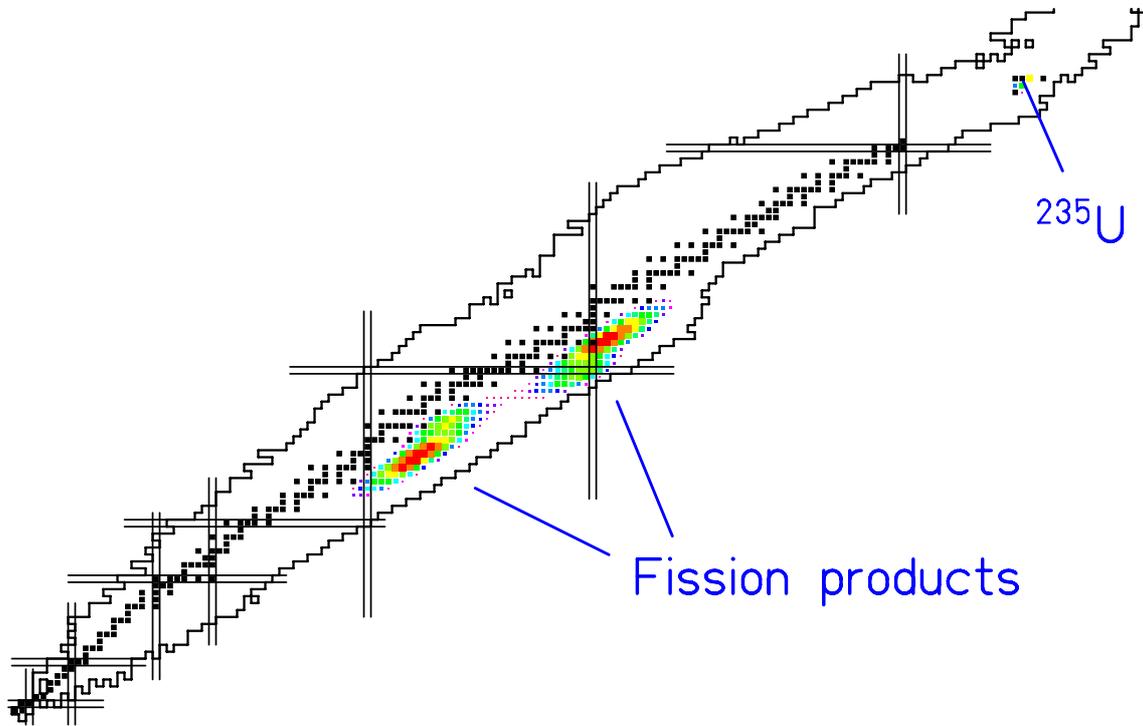
Fission



- **Basic characteristics:**
 - Curvature of stability valley due to Coulomb repulsion.
 - **Fission products are neutron-rich.**
 - Shell effects (^{132}Sn) or fluctuations in polarisation.
 - **The only ways to reach even more neutron-rich.**
 - The asymmetry term in the liquid drop is large.
 - **Polarisation is small.**
- **Excitation energy leads to opposite effects**
 - Fluctuations in polarisation increase.
 - **Fission extends to more neutron-rich isotopes.**
 - Excited fission fragments evaporate neutrons.
 - **Final fission products are less neutron-rich.**

Fission of actinides

Fission induced by low-energy neutrons

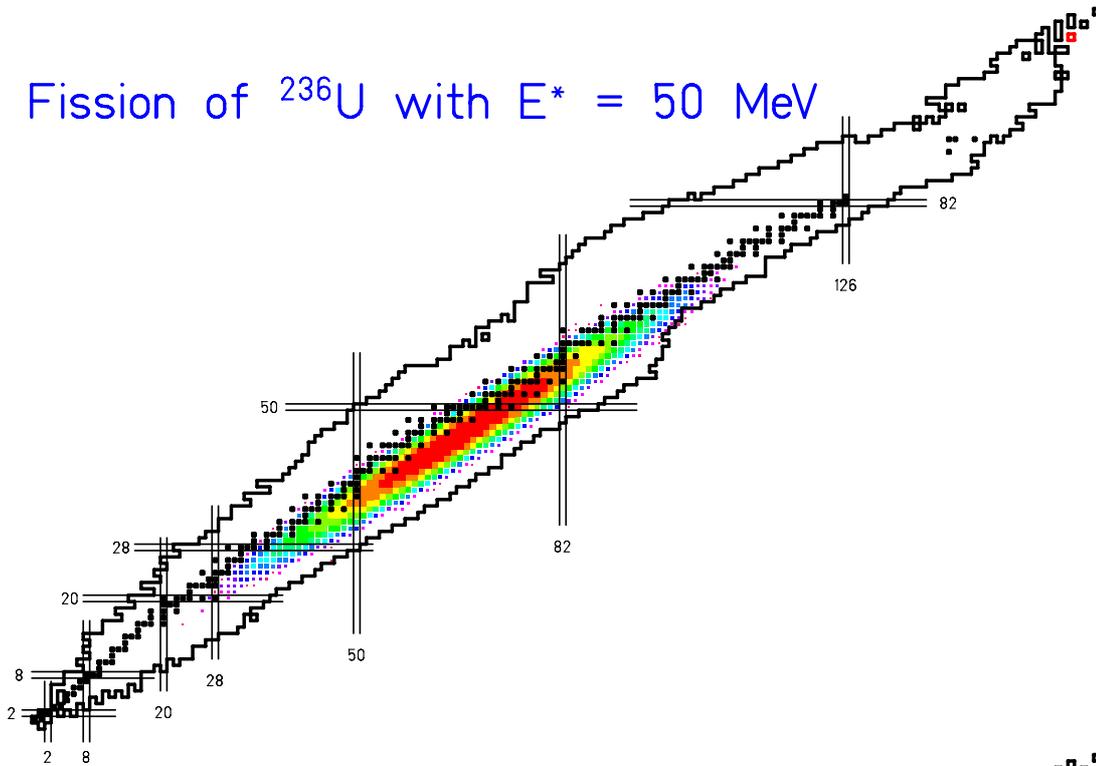


Fission of actinides from low excitation energies induced by neutrons, protons, electrons, photons.

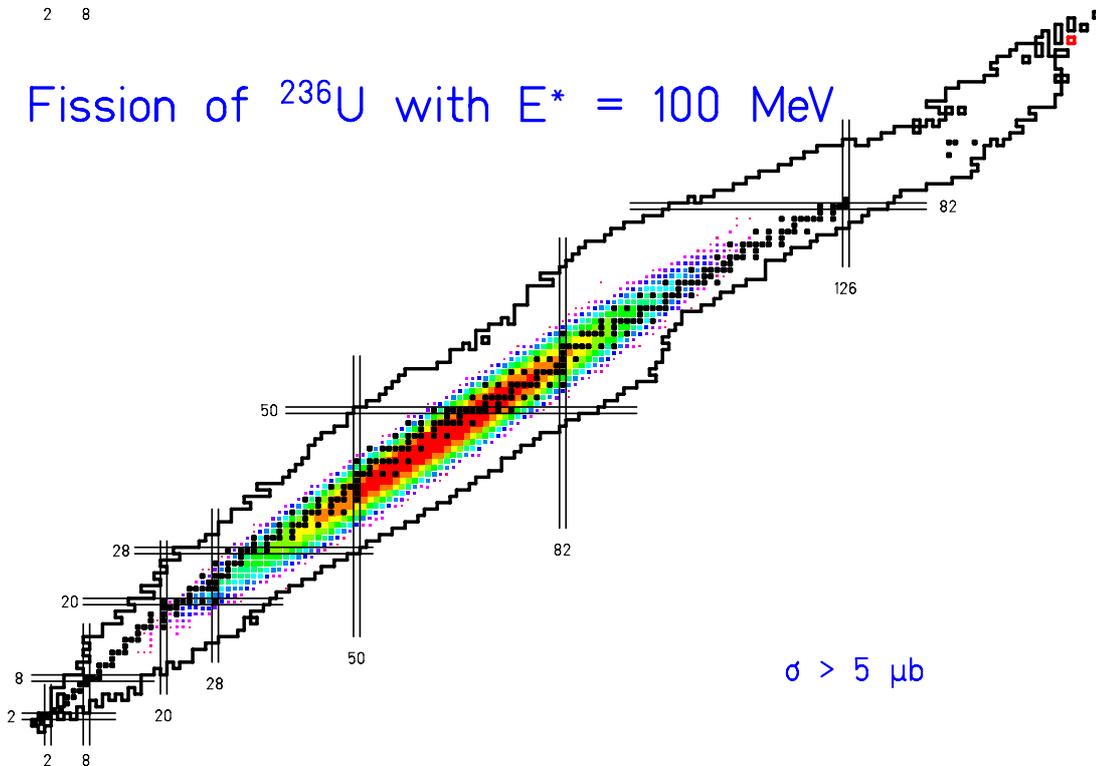
Production of moderately neutron-rich isotopes of a few elements.

Fission: Variation of Excitation Energy (Calculations)

Fission of ^{236}U with $E^* = 50$ MeV



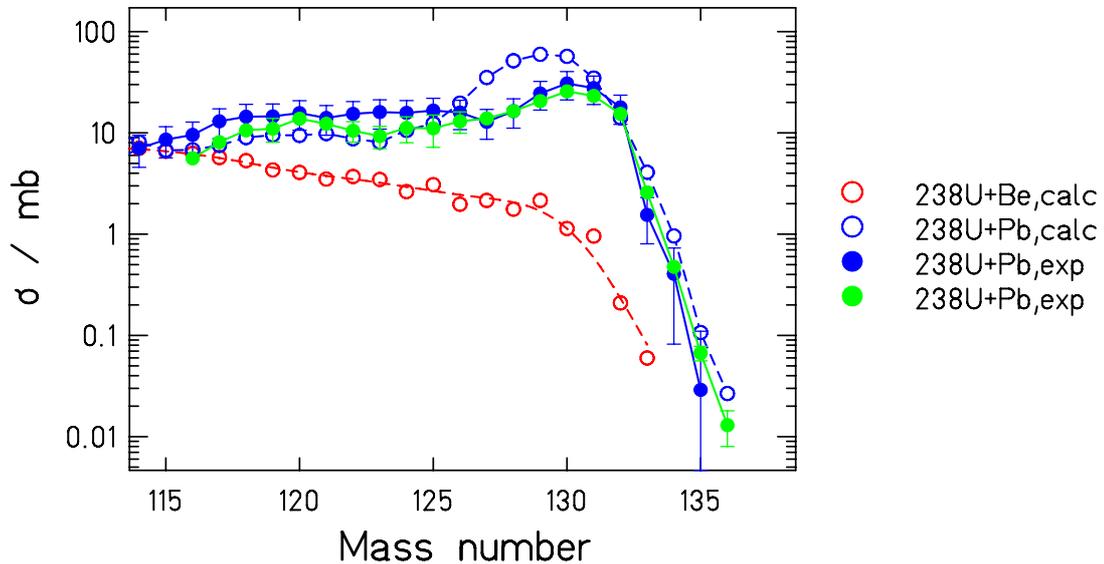
Fission of ^{236}U with $E^* = 100$ MeV



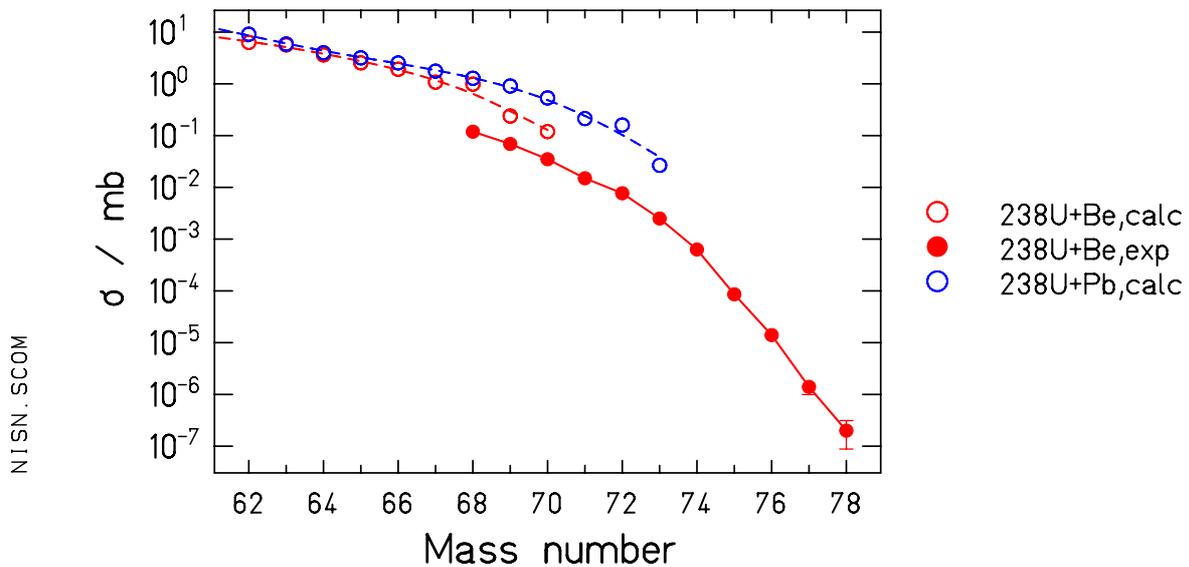
Higher E^* → larger fluctuations in A
→ less neutron-rich

Exploring the limits of neutron-rich isotopes by fragmentation-fission reactions

Production of Sn isotopes by fission



Production of Ni isotopes by fission

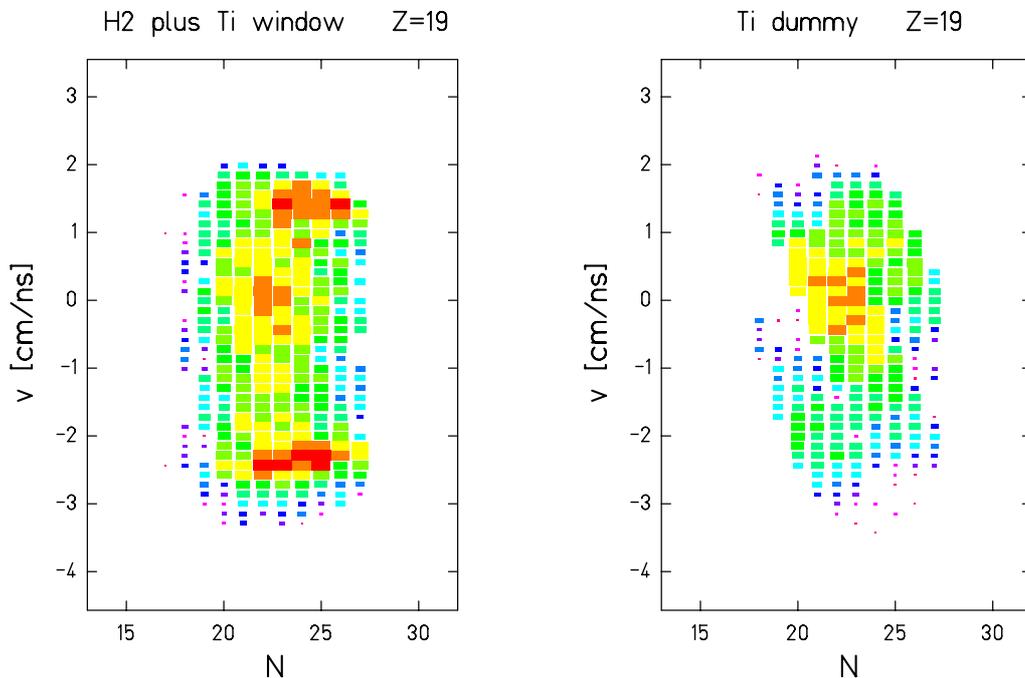


NISN . SCOM

Data from T. Enqvist et al., Nucl. Phys. A 658 (1999) 47,
C. Donzaud et al., Eur. Phys. J. A1 (1998) 407,
C. Engelmann et al., Z. Phys. A 352 (1995) 351.

**Step decrease of cross sections due to limitation of charge
polarisation in fission.**

Kinematic Properties of Potassium Produced from ^{238}U in Different Targets



Projectile: ^{238}U , 1 A GeV

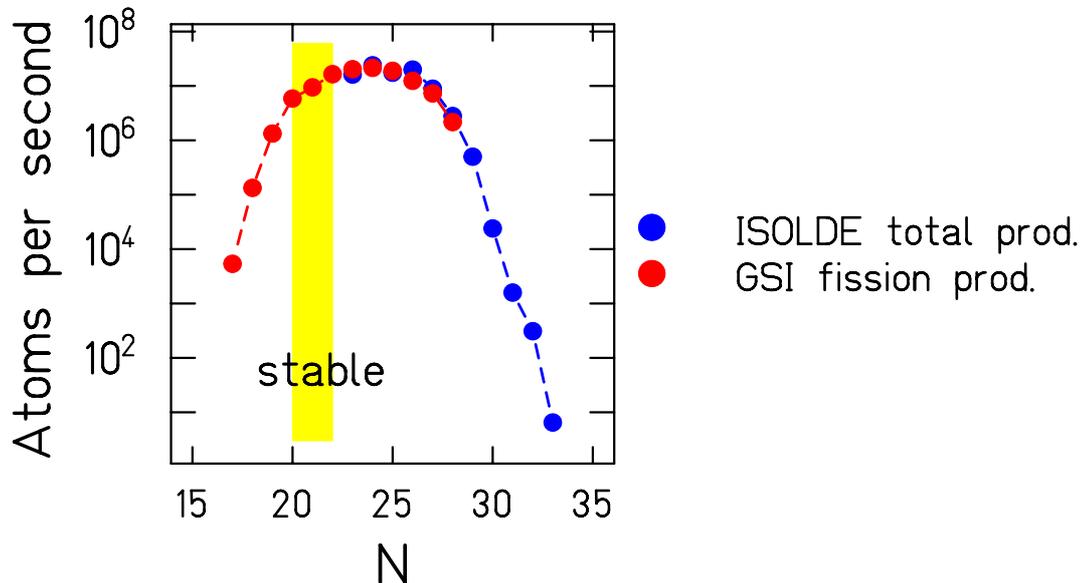
Target left: hydrogen (+ titanium window)
right: titanium

Velocity distributions of potassium isotopes

- Production in hydrogen target from very asymmetric fission.
- Production in titanium target from projectile fragmentation.

Data from M. V. Ricciardi, GSI, thesis in preparation.

Production of Potassium in $p + {}^{238}\text{U}$



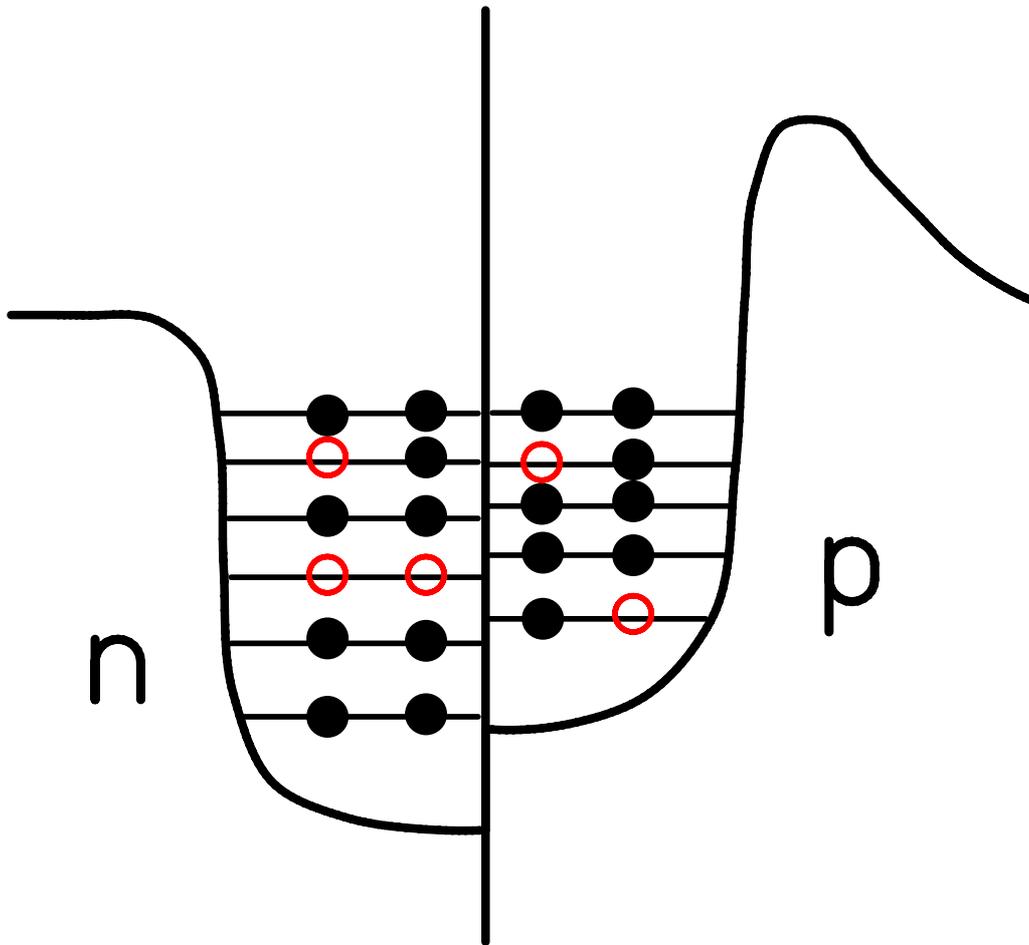
Isotopic yields from 600 MeV protons on ${}^{238}\text{U}$ (ISOLDE)
and
fission-product yields from 1 A GeV ${}^{238}\text{U}$ + hydrogen (GSI)

No absolute cross sections from ISOLDE yields.

The distributions fit together:
ISOLDE yields of light elements from fission!

Data from
H.-J. Kluge, ISOLDE user's guide, CERN 86-05 (1986)
and
M. V. Ricciardi, GSI, thesis in preparation

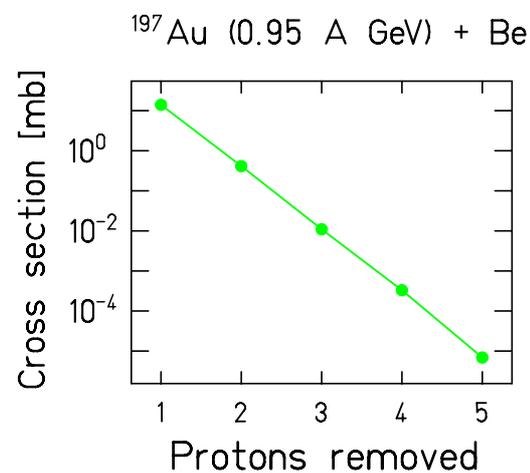
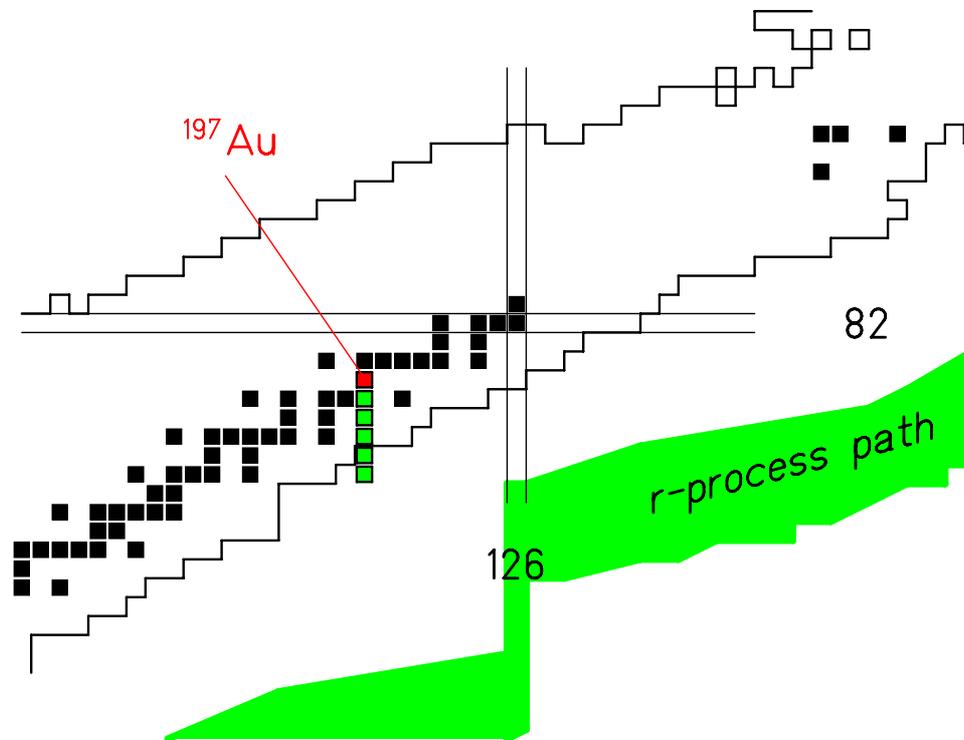
Fragmentation



Removal of nucleons in quasi-free nucleon-nucleon collisions.

- **Large fluctuations in N/Z .**
- **Large fluctuations in excitation energy.**

Dedicated study of proton-removal channels



Data from J. Benlliure et al., Nucl. Phys. A 660 (1999) 87.

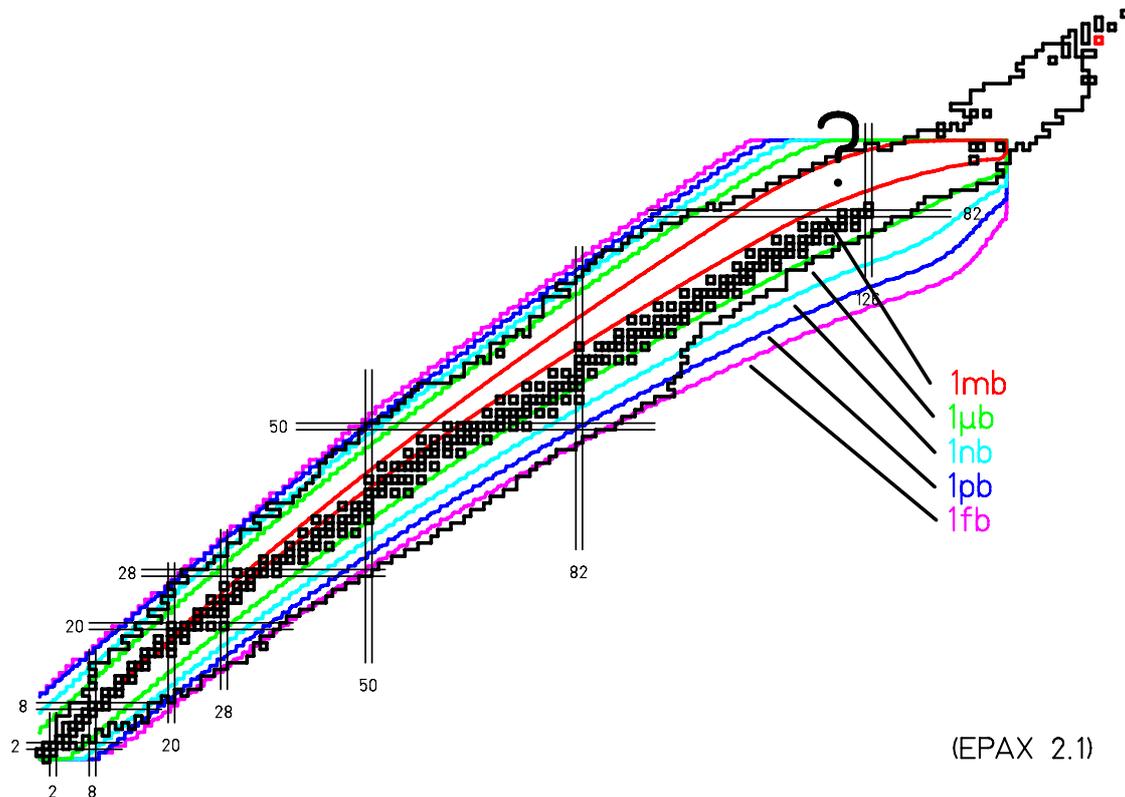
Abrasion of protons only, no evaporation of neutrons

=> "Cold Fragmentation".

Promising results for producing neutron-rich nuclei.

Expected production cross sections by cold fragmentation

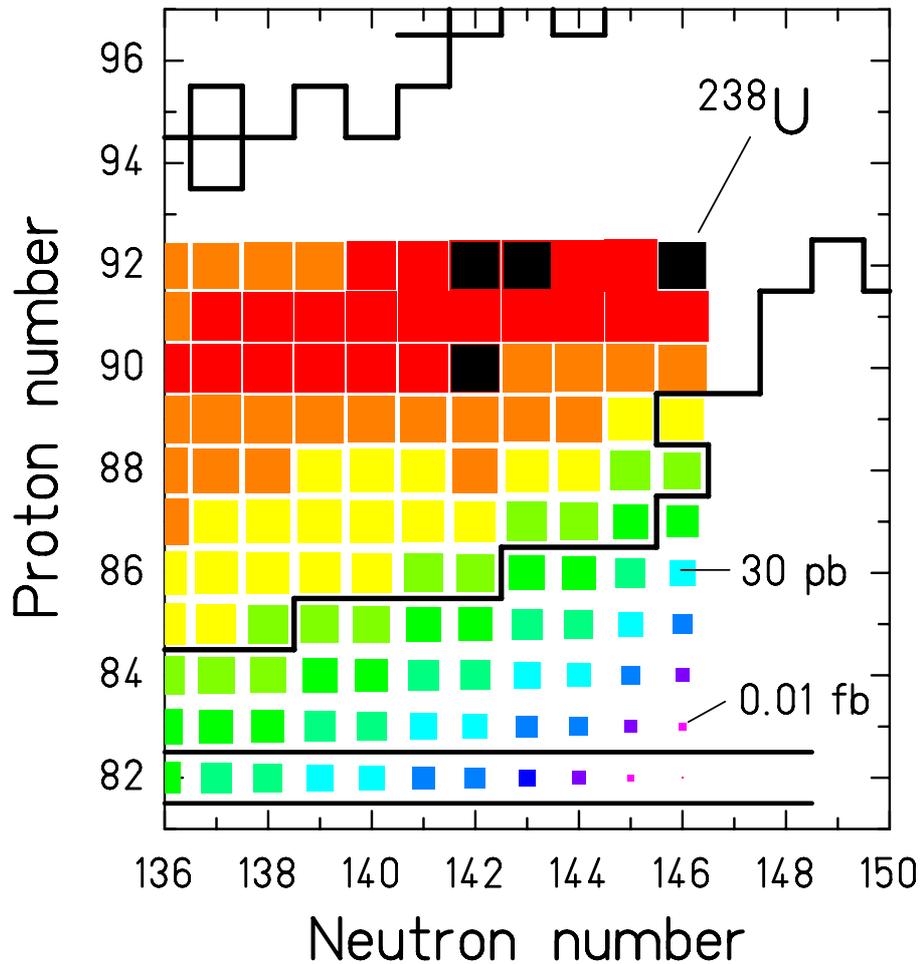
Isotopic production cross sections, $^{238}\text{U} + ^7\text{Be}$



The empirical systematics EPAX which has carefully been adjusted to available experimental data has been used to estimate isotopic production cross sections of extremely neutron-rich isotopes by cold fragmentation. Since EPAX does not consider fission, the prediction of neutron-deficient isotopes is not realistic.

K. Sümmerer, B. Blank, Phys. Rev. C 61 (2000) 034607

Reaching extremely neutron-rich isotopes by cold fragmentation



Result of a model calculation^{a)} on cold fragmentation ^{238}U (1 A GeV) + Be. Many new neutron-rich isotopes in reach, but low cross sections!

^{a)} The model is described in J. Benlliure et al., Nucl. Phys. A 660 (1999) 87